

**Dislocation Dynamics in Metals at Atomic-scale:  
Interactions between Dislocations and Obstacles with Dislocation Character**

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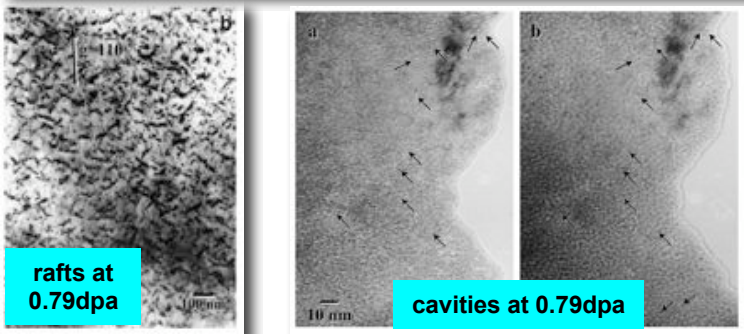
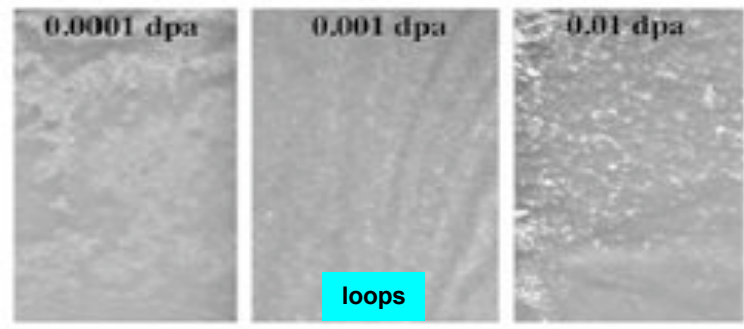
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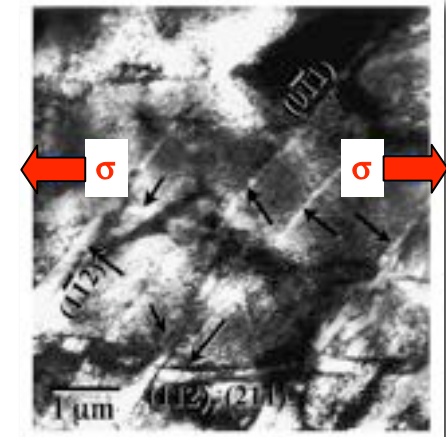
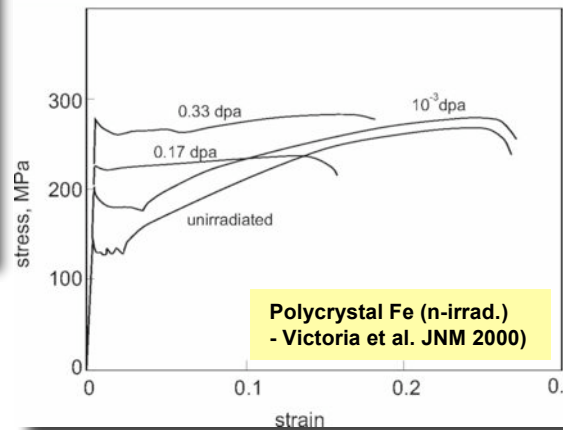
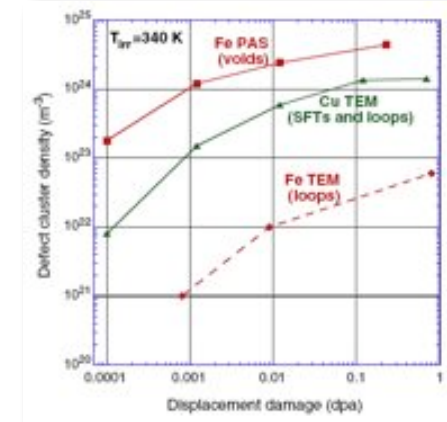
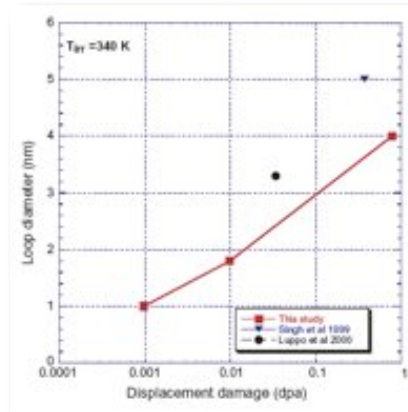
In collaboration with:

Dmitry Terentyev (CEN/SCK, Belgium)

## Motivation: effect of irradiation microstructure on mechanical properties



Fe, n-irradiated at 60°C (Zinkle & Singh JNM 2006)

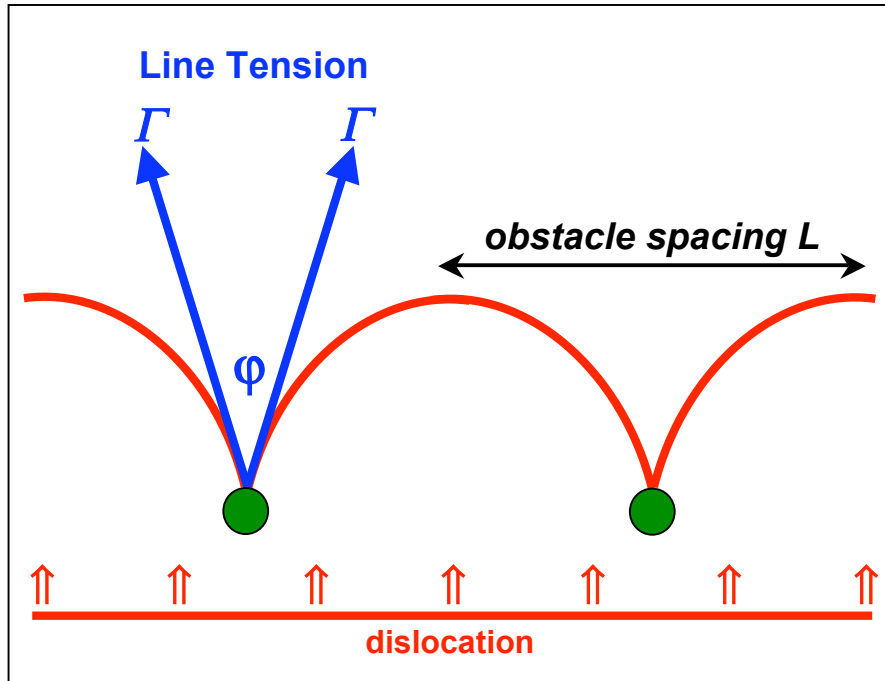


Deformed Fe, n-irrad. 0.4 dpa (Zinkle & Singh JNM 2006)

- dislocations under stress move through field of irradiation-induced obstacles
  - dislocation loops, SFTs, point defect clusters, voids, precipitates, etc.
- yield and flow stress raised, strain to failure reduced

## Continuum modelling of strengthening

### (a) Line tension approximation



- Dislocation overcomes obstacle when  $\varphi = \varphi_c$  at  $\tau = \tau_c$

$$\tau_c b L = 2 \Gamma \cos(\varphi_c / 2)$$

- $\Gamma \approx G b^2 / 2$

$$\therefore \tau_c = \alpha G b / L$$

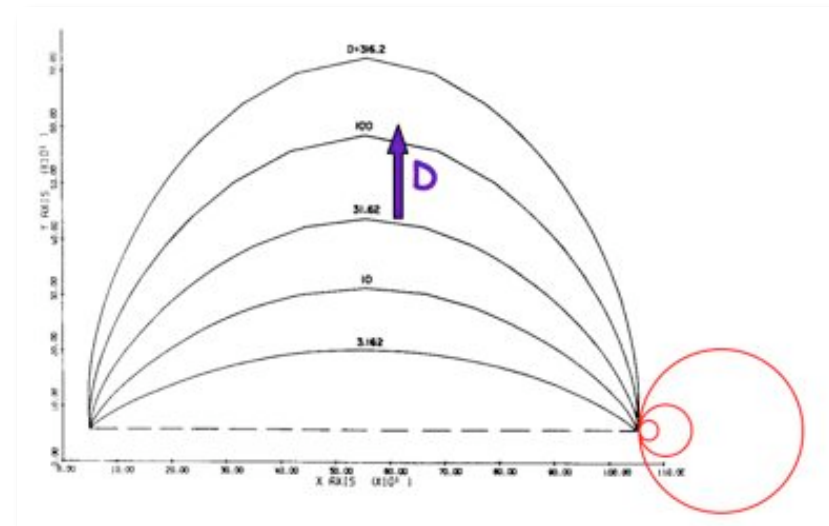
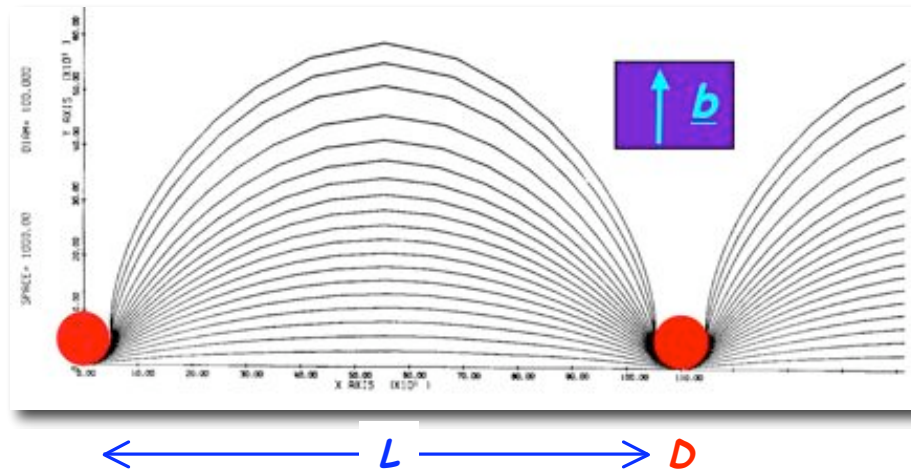
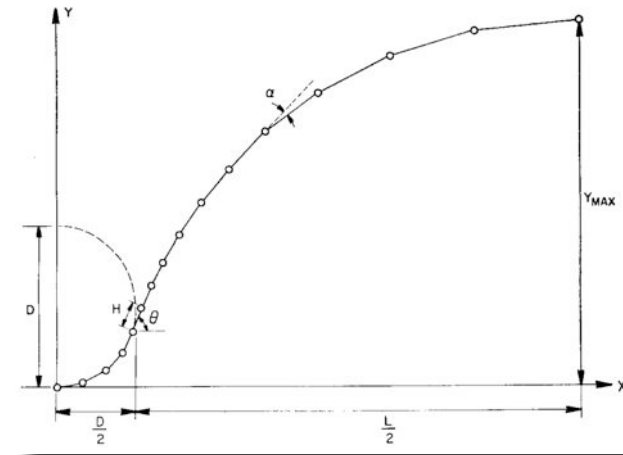
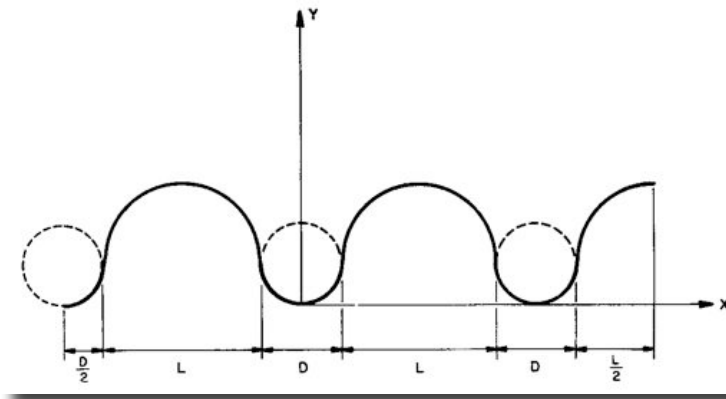
$$\alpha = \cos(\varphi_c / 2)$$

- characterises 'obstacle strength'

- $\varphi_c$  is largely empirical
- strong obstacles  $\varphi_c \sim 0$ , but empirically  $\alpha \sim 0.2-0.5$
- LT model ignores dislocation self-interaction
- self-stress is included in 'dislocation dynamics' (DD) modelling (elasticity approximation)

## (b) Dislocation self-stress simulation for strong obstacles

- Orowan strengthening (impenetrable obstacles)



[Bacon, Scattergood, Kocks (Phil. Mag. 1973-)]

- critical stress for edge dislocation and impenetrable particles or voids:

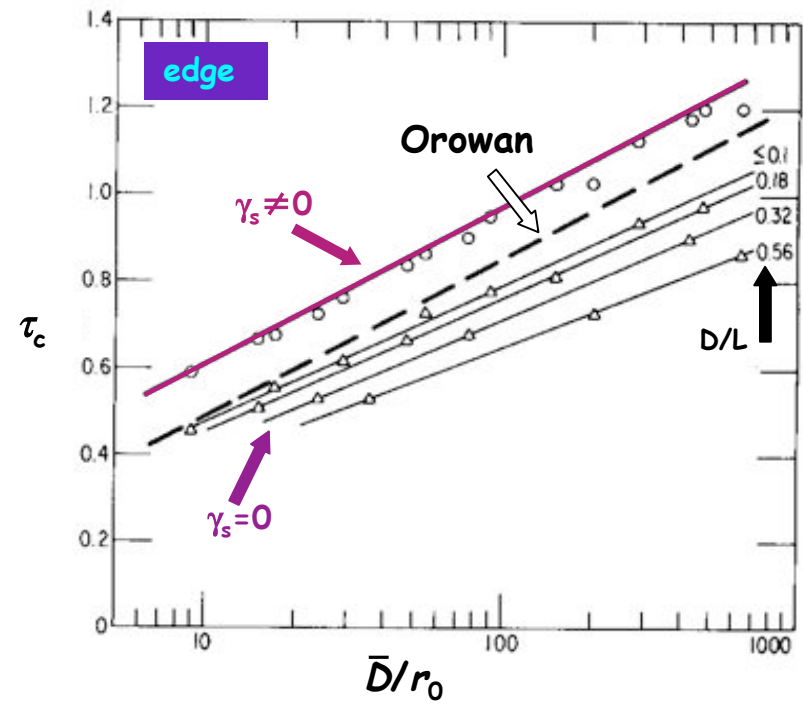
$$\tau_c \approx \frac{Gb}{2\pi L} \left[ \ln \left( \frac{1}{D^{-1} + L^{-1}} \right) + B \right]$$

where  $B$  depends on  $\gamma_s$  for voids

- cf line tension:  $\tau_c = \alpha Gb/L$
- energy (tension) of dipole  $\propto \ln[D]$   
 $\Rightarrow$  size-dependence for strong obstacles
- dislocation dynamics (DD) requires 'local rules'
  - effects of core structure
  - mobility of segments, strength of junctions
  - dependence on  $\tau$ ,  $T$ ,  $\dot{\epsilon}$

⇓ link to atomic scale

- computer simulation of atomic-scale processes by which obstacles affect dislocation motion  
 $\Rightarrow$  quantitative data on  $\tau$ ,  $T$ ,  $\dot{\epsilon}$



## (c) Atomic-scale simulations of dislocation-obstacle interaction

- System size
  - few  $M$  mobile atoms
- Obstacle
  - periodic spacing  $L \leq 100\text{nm}$
  - size  $D \leq 10\text{nm}$
  - $\rho_D \sim 10^{14}-10^{15}\text{m}^{-2}$

- Statics (MS)  $T = 0\text{K}$

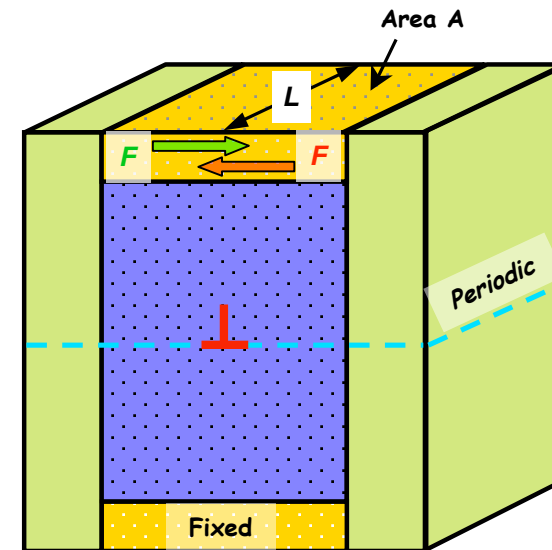
- apply  $\Delta\varepsilon$  incrementally
- relax to minimum pot'l energy
- equivalent to elasticity

- Dynamics (MD)  $T > 0\text{K}$

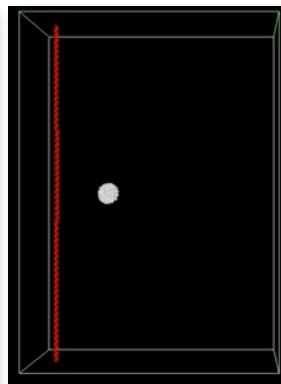
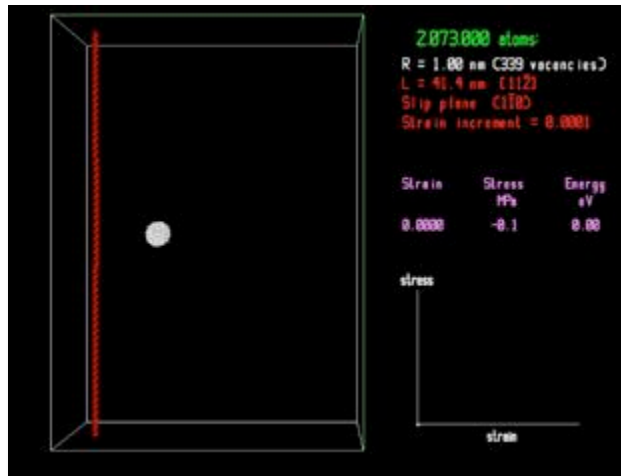
- either apply strain rate  $\sim 10^6-10^8\text{s}^{-1}$
- dislocation veloc  $\sim 5-500\text{ms}^{-1}$
- or apply stress

$$\tau_{\text{appl}} = -F/A$$

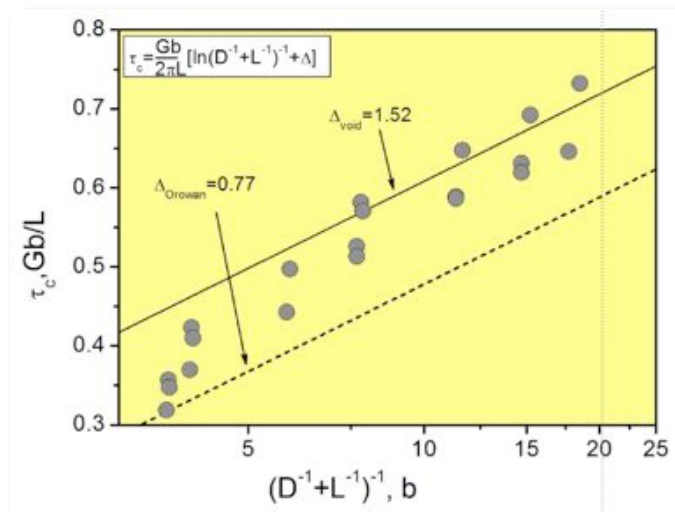
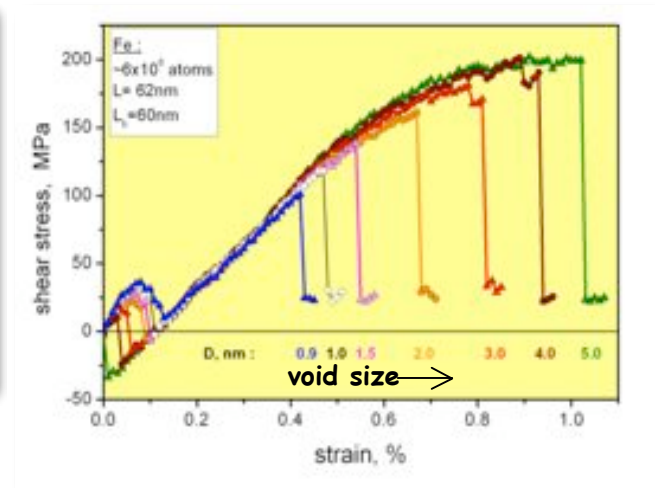
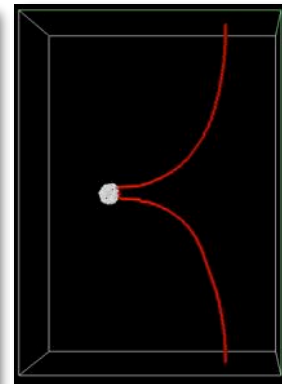
$$\tau_{\text{appl}} = F/A$$



## Void strengthening $\alpha$ -Fe at $T = 0\text{K}$ (edge dislocation)



2nm (339v) void in Fe  
 L=42nm, T=0K



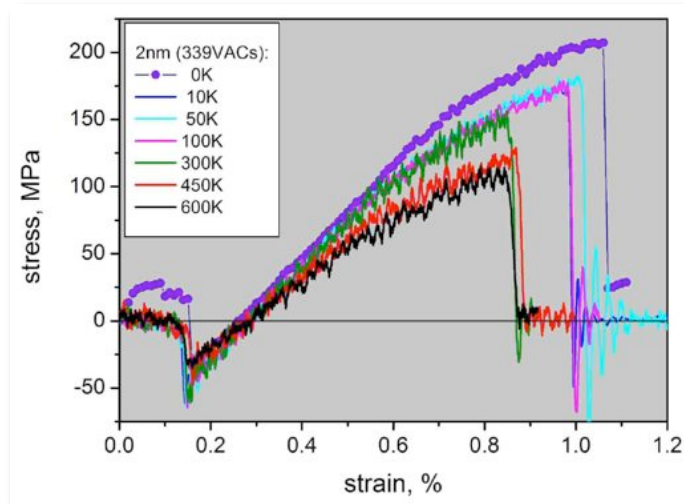
$$\tau_c \approx \frac{Gb}{2\pi L} \left[ \ln \left( \frac{1}{D^{-1} + L^{-1}} \right) + B \right]$$

- similar dependence on  $L$  and  $D$  from atomic-level and continuum treatments

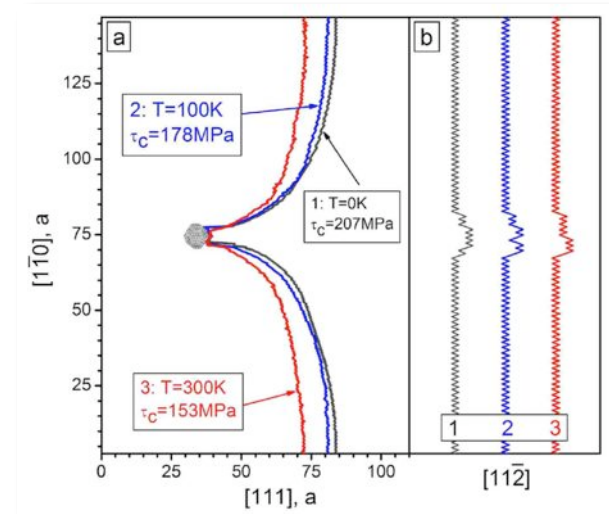
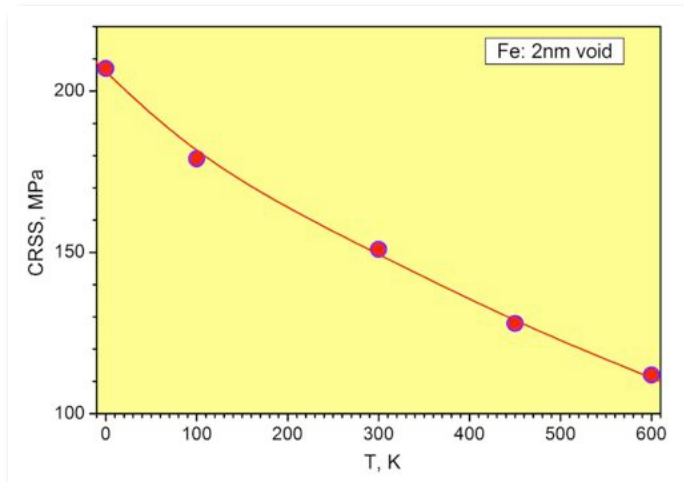


## Void strengthening at $T > 0K$ (edge dislocation)

- motion under constant strain-rate ( $10^6$ - $10^8s^{-1}$ ) - dislocation dynamics at the atomic scale

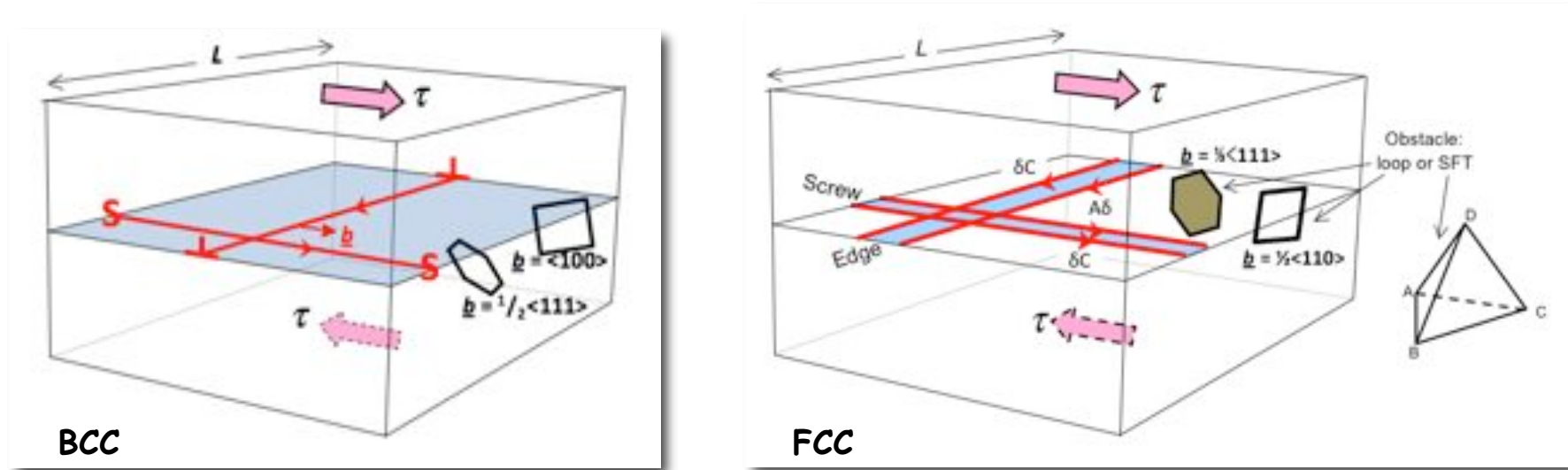


e.g.  $\tau_c$ : 2nm voids in Fe  
- strain rate  $5 \times 10^6 s^{-1}$





(d) Dislocation interaction with nano-scale dislocation obstacles



- Only consider BCC here

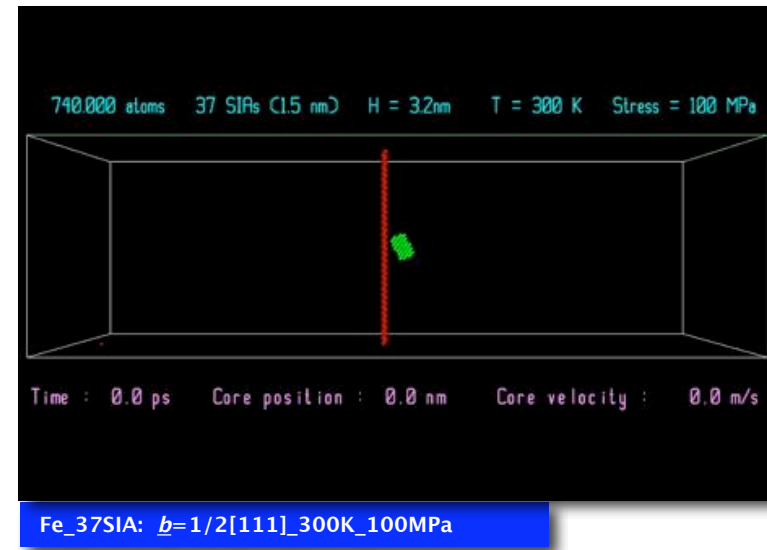
• Outcomes

- R1: dislocation and obstacle unchanged - edge or screw
- R2: obstacle changed but dislocation unchanged - edge or screw
- R3: partial or full absorption of obstacle by edge dislocation (superjog formation)
- R4: temporary absorption of obstacle by screw dislocation (helix formation)
- loop drag ( $\equiv$  R1+R3)

[Bacon, Osetsky & Rodney, in *Dislocations in Solids*, vol. 15 (2009)]

**No intersection:** Small  $\frac{1}{2}[111]$  loop in Fe at 300K

- edge dislocation
- parallel bs
- loop drag
- breakaway above critical stress



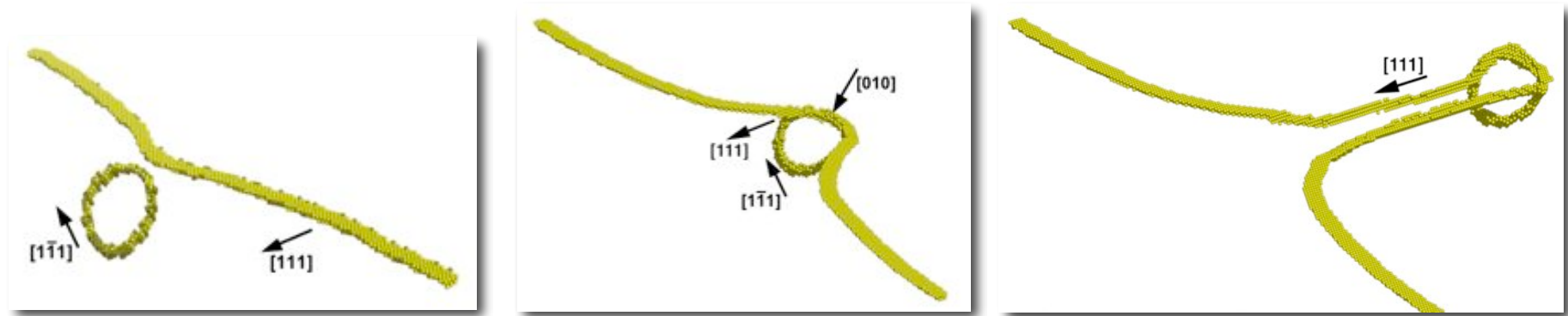
[Rong, Osetsky & Bacon, Phil. Mag. (2005)]

**Intersection:** Frank's rule for energetically-favourable reactions in Fe (BCC) predicts

- $\frac{1}{2}\langle 111 \rangle$  loops acquire  $\langle 100 \rangle$  segments
  - e.g.  $\frac{1}{2}[111] + \frac{1}{2}[\bar{1}\bar{1}\bar{1}] = [100]$
- $\langle 100 \rangle$  loops acquire  $\frac{1}{2}\langle 111 \rangle$  segments
  - e.g.  $\frac{1}{2}[111] + [\bar{1}00] = \frac{1}{2}[\bar{1}11]$

## R1: dislocation and obstacle unchanged

**Intersection:** Large loop (331i) with  $\underline{b} = \frac{1}{2}[1-11]$  in Fe at 100K  
- edge dislocation  $\underline{b} = \frac{1}{2}[111]$



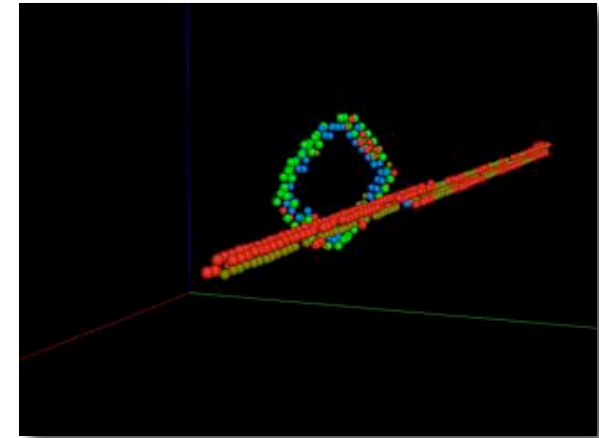
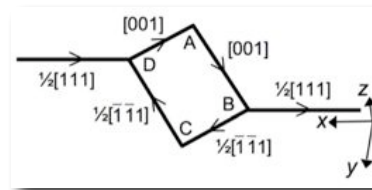
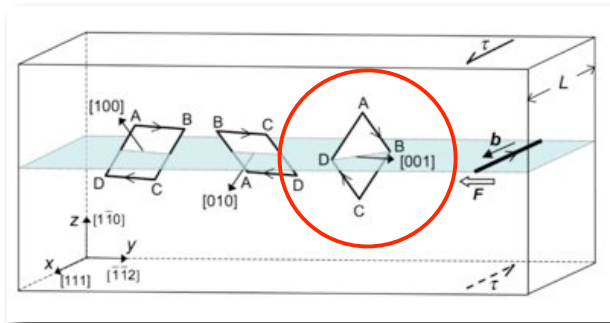
- attractive reaction forms [010] segment
- [010] segment has different glide plane and is immobile at low  $T$
- screw dipole drawn out
- pinches off by  $\alpha$ -slip at  $\tau_c$ , leaving  $\frac{1}{2}[1-11]$  loop

⇒ strong obstacle

[Bacon, Osetsky & Rong, Phil. Mag. (2006)]

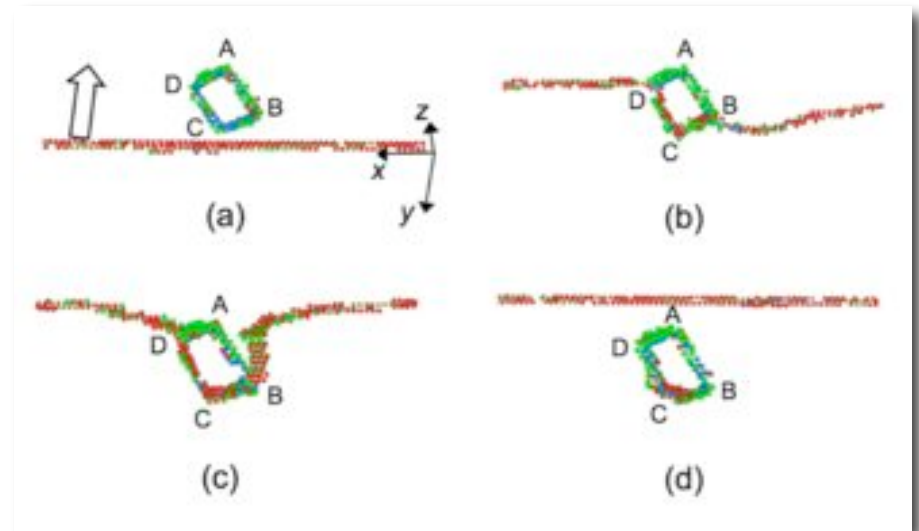
## R1: dislocation and obstacle unchanged (cont'd)

**Intersection:** Large loop (128i) with  $\underline{b} = [001]$  in Fe at 300K  
 - screw dislocation  $\underline{b} = \frac{1}{2}[111]$



- attractive reaction converts sides BC & CD to  $\frac{1}{2}[-1-11]$  by x-slip of screw
- screw side arms x-slip to corner C at  $\tau_c$  leaving original loop
- screw glide plane now coincident with C (periodic boundaries)

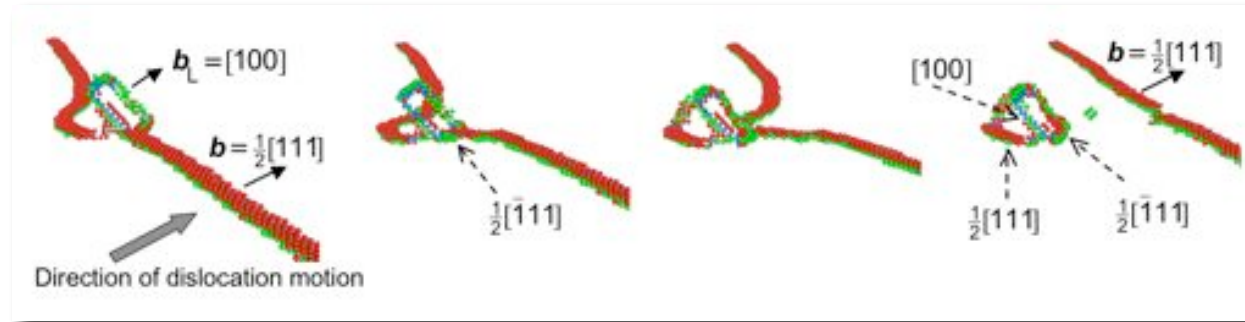
⇒ moderately strong obstacle



[Terentyev, Bacon & Osetsky, Phil Mag, in press]

## R2: dislocation unchanged, obstacle changed

Intersection:  $[100]$  loop (169i) in Fe at 300K  
 -  $\frac{1}{2}[111]$  edge dislocation

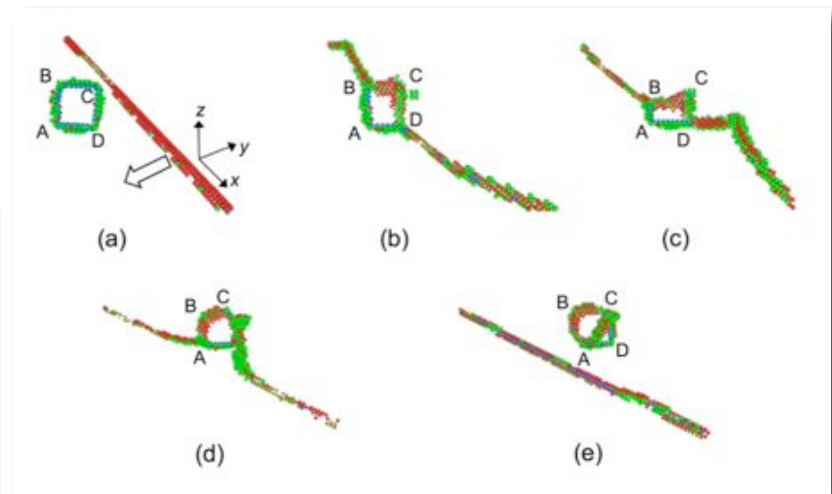
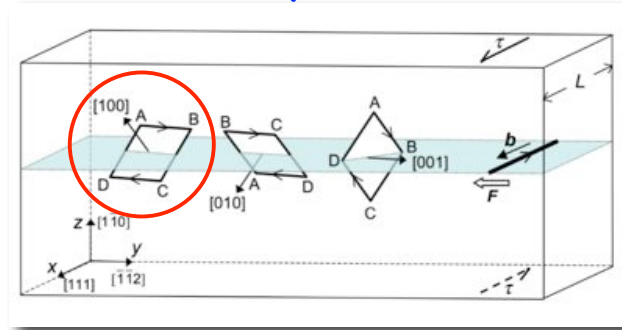


- dislocation repelled but forms  $\frac{1}{2}[-111]$  segment on contact
- double loop complex remains

[Terentyev, Bacon & Osetsky, Acta Mat. (2008)]

⇒ strong obstacle

- similar complex for  $[100]$  loop and  $\frac{1}{2}[111]$  screw:

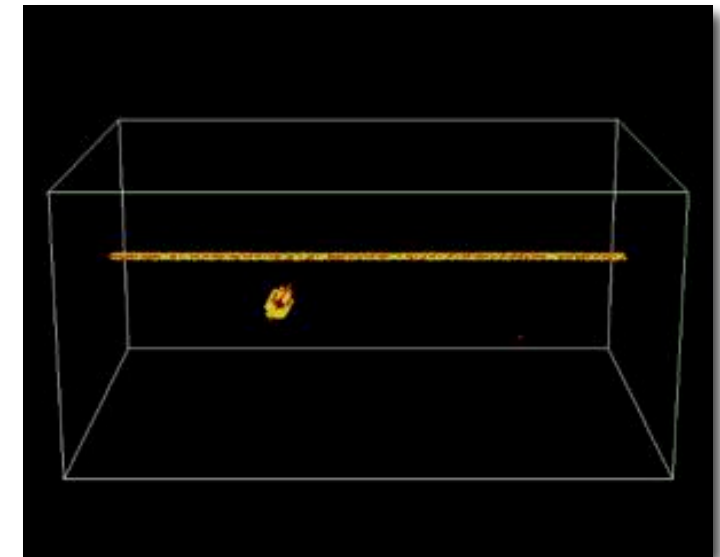
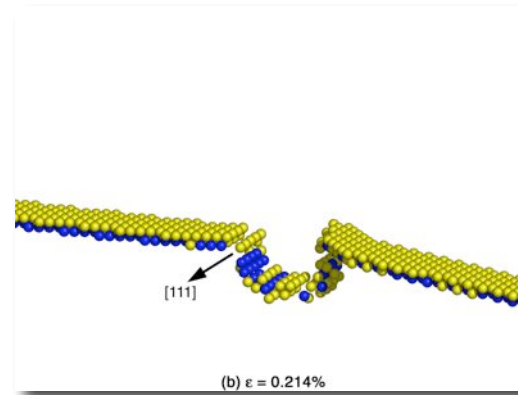
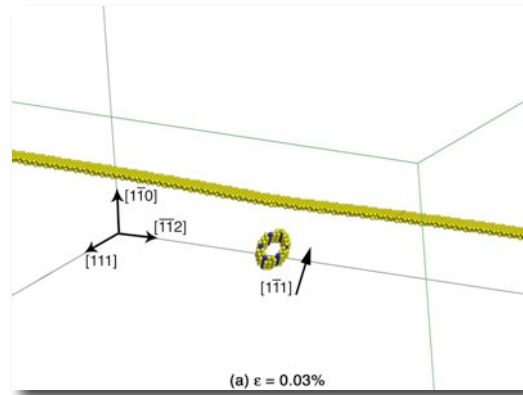


[Terentyev, Bacon & Osetsky, Phil Mag, in press]

### R3: partial or full absorption of obstacle by edge dislocation (superjog formation)

Small loop (37i) with  $\underline{b} = \frac{1}{2}[1-11]$  in Fe at 0-600K

- edge dislocation  $\underline{b} = \frac{1}{2}[111]$



Fe\_37SIA:  $\underline{b}=1/2[1-11]$ \_300K\_SR10

[Bacon, Osetsky & Rong, Phil. Mag. (2006)]

• loop  $\underline{b}$  changes on contact

⇒ weak obstacle

• efficient absorption of all SIAs



## R3: partial or full absorption of obstacle by edge dislocation (superjog form) (cont'd)

Large loop (331i) with  $\underline{b} = \frac{1}{2}[1-11]$  in Fe at 300-600K

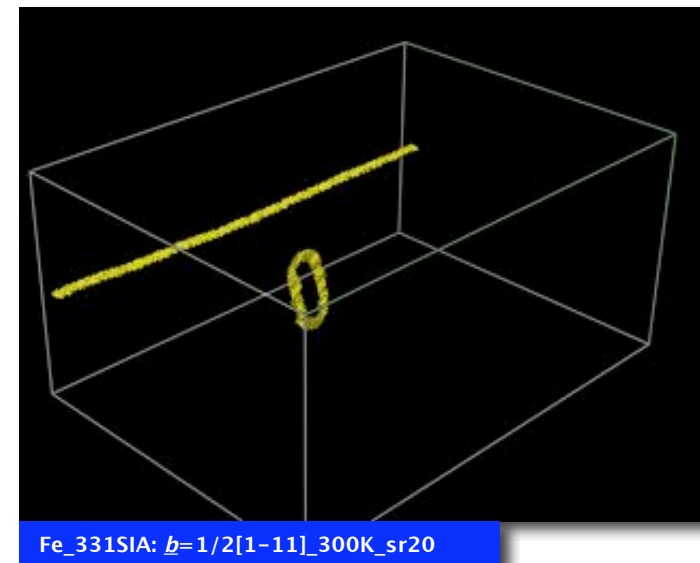
- edge dislocation  $\underline{b} = \frac{1}{2}[111]$



- sessile [010] segment forms on contact
- segment has low mobility
- glides over loop converting it to  $\frac{1}{2}[111]$  as screw side arms x-slip

⇒ strong obstacle

- efficient absorption of all SIAs



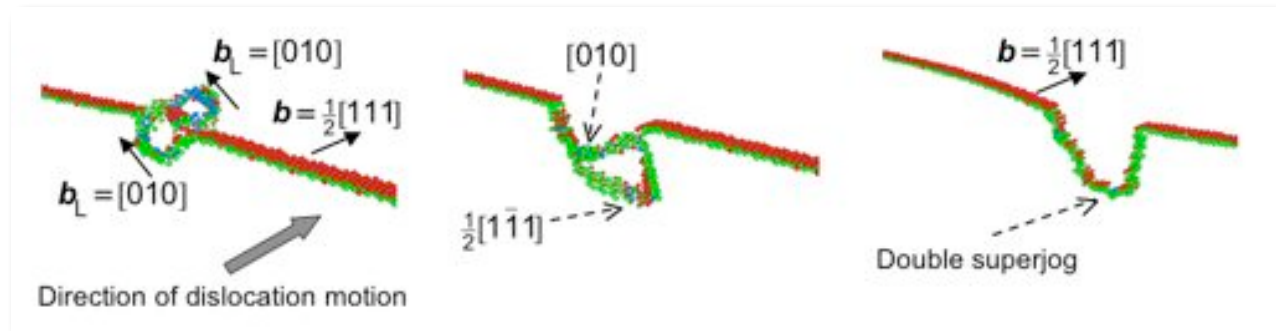
[Bacon, Osetsky & Rong, Phil. Mag. (2006)]



### R3: partial or full absorption of obstacle by edge dislocation (superjog form<sup>n</sup>) (cont'd)

[010] loop (169i) in Fe at 300K

-  $\frac{1}{2}[111]$  edge dislocation



- $\frac{1}{2}[1-11]$  segment forms on contact
- segment glides down  $\Rightarrow$  same configuration as interaction with  $\frac{1}{2}[1-11]$  loop above
- segment glides over loop converting it to  $\frac{1}{2}[111]$

$\Rightarrow$  strong obstacle

- efficient absorption of all SIAs

[Terentyev, Bacon & Osetsky, Acta Mat. (2008)]

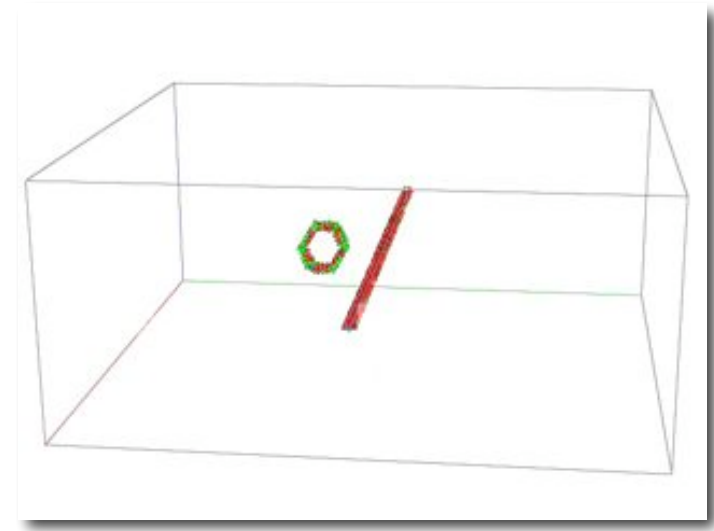
**R4: temporary absorption of obstacle by screw dislocation (helical turn formation)**  
- followed by detachment as turn closes

$\frac{1}{2}[111]$  loop in Fe at 100K

- screw dislocation  $\underline{b} = \frac{1}{2}[111]$

- loop absorbed as helical turn
- cannot glide with line
- line released when turn closes and loop restored

⇒ strong obstacle



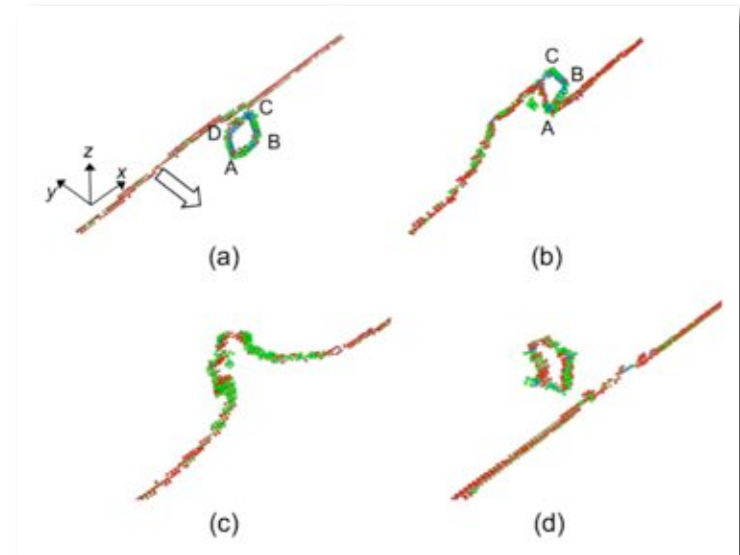
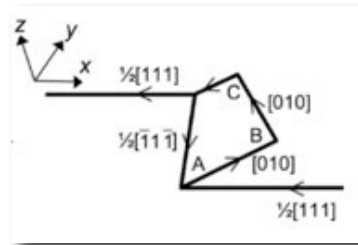
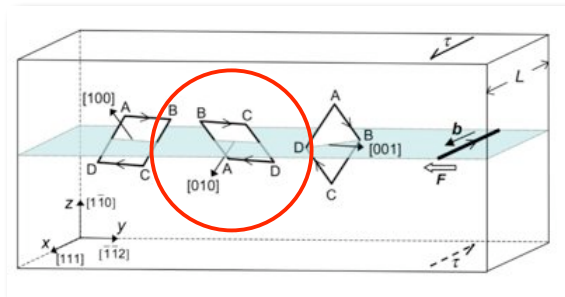
[Terentyev, unpublished]

- any net absorption/transport of SIAs is along the line

**R4: temporary absorption of obstacle by screw dislocation (helical turn formation)**  
 - followed by detachment as turn closes (cont'd)

[010] loop in Fe at 300K

- screw dislocation  $b = \frac{1}{2}[111]$



- screw initially repelled, but x-slips to corner D and converts AD to  $\frac{1}{2}[1-11]$  segment
- segment sweeps over loop as screw side arms cross-slip, converting other sides to  $\frac{1}{2}[111]$   
 $\Rightarrow$  loop absorbed as helical turn on screw
- $\frac{1}{2}[111]$  loop formed when screw breaks away

$\Rightarrow$  strong obstacle

- no net absorption/transport of SIAs

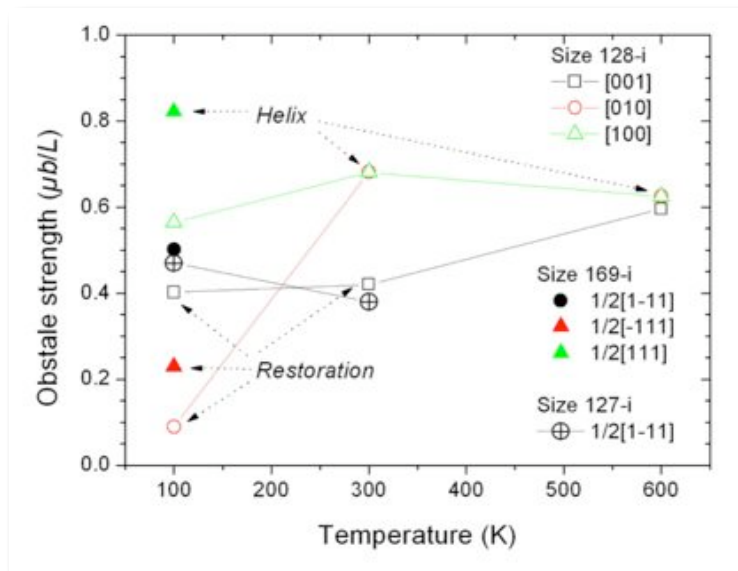
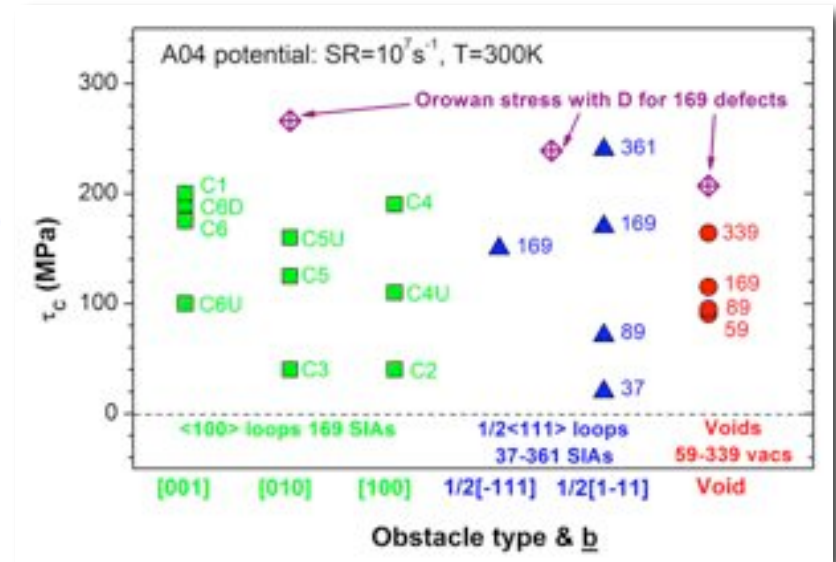


[Terentyev, Bacon & Osetsky, Phil Mag, in press]

## Summary/conclusions for $\tau_c$ in Fe

$\tau_c$  at 300K for edge with  $L = 41\text{nm}$

- large variation in  $\tau_c$
- large variation in defect absorption on line (0-100%)
- no correlation between  $\tau_c$  and absorption



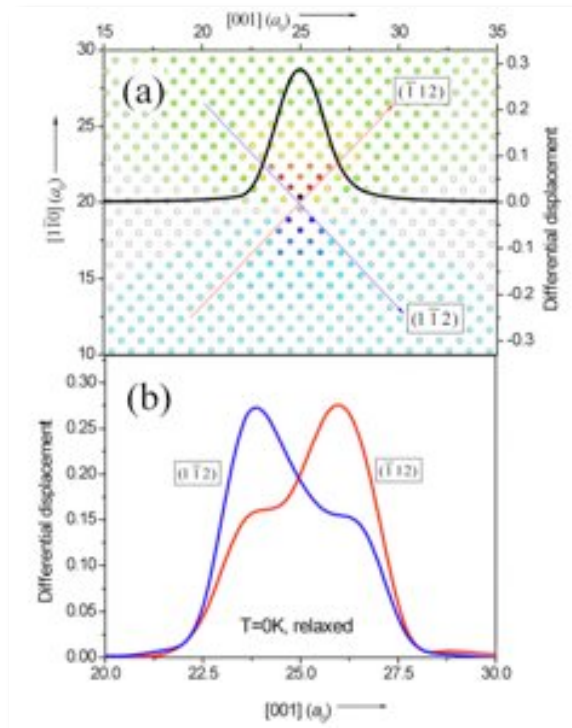
$\tau_c$  for screw at 100-600K for  $L = 25\text{nm}$

- reactions
  - R1 (loop restored)  $\rightarrow$  lowest  $\tau_c$
  - R2 (loop changed)  $\rightarrow$  medium  $\tau_c$
  - R4 (loop absorbed in temporary helix)  $\rightarrow$  highest  $\tau_c$

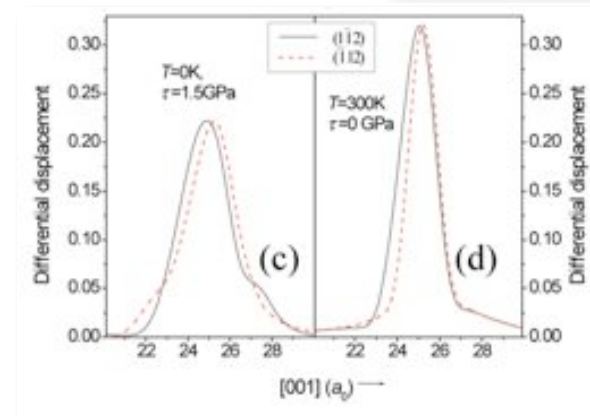
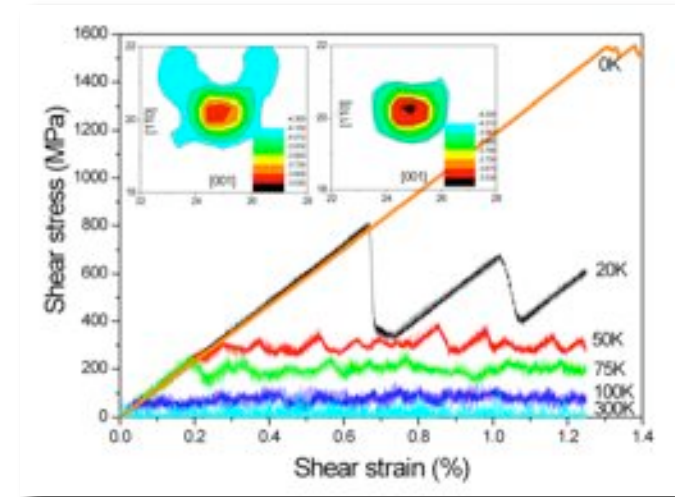
[Terentyev, Bacon & Osetsky, Phil Mag, in press; Liu & Biner, Scripta Mat. 2008]

## What else?

- all reactions to form new segments, e.g.  $\langle 100 \rangle$  on  $\frac{1}{2}\langle 111 \rangle$  loop or  $\frac{1}{2}\langle 111 \rangle$  on  $\langle 100 \rangle$  loop, satisfy Frank's Rule:  $(b_1^2 + b_2^2) > b_3^2 \iff$  even for nm loops
- final outcome depends on mobility of  $\langle 100 \rangle$  dislocation, and therefore on  $T$ :



3-D core of  $\langle 001 \rangle$  edge:  
 $\underline{b} = \frac{1}{2}\langle 111 \rangle$  on two  $\{112\}$  planes



3-D core  $\Rightarrow$  2-D at high  $\tau$  and/or  $T$   
 $\underline{b} = \frac{1}{2}[001]$  on two  $\{112\}$  planes

[Terentyev, Osetsky & Bacon, Acta Mater. in the press]

## Conclusions

### Dislocation Dynamics in Metals at Atomic-scale: Interactions between Dislocations and Obstacles with Dislocation Character

- need for predictive computer modelling
- atomic scale simulation can provide unique information on mechanisms
  - qualitative and quantitative
  - wide variety of nano-scale obstacles
  - can validate continuum models
  - increasing understanding of reactions, outcomes, obstacle strength
- future challenges
  - more realistic interatomic potentials, e.g. Fe, alloys
  - multiple dislocation effects, e.g. channelling
  - grain and interphase boundaries
  - strain rate effects
  - local rules and activation parameters for continuum-based dislocation dynamics